# On the Extrapolation of Band-Limited Functions with Energy Constraints 

WEN YUAN XU and CHRISTODOULOS CHAMZAS


#### Abstract

A modification of the algorithm proposed by Papoulis and Gerchberg for extrapolating band-limited functions ${ }^{1}$ is suggested to extend its applicability over data corrupted by noise. We assume that energy constraints are known either for the band-limited signal or for the noise. In addition, the discrete formulation of the iterative algorithm is derived, and the transition from the continuous algorithm to its digital implementation is presented.


## I. Introduction

THE extrapolation of a band-limited function $f(t)$ for any $t$ in terms of a finite segment $g(t)$ of $f(t)$ is an essential problem in signal analysis. Various algorithms have been proposed for this problem. All of them can, essentially, be classified in three categories: algorithms based on analytic continuation, algorithms using a series expansion in terms of prolate spheroidal functions [1], [7] (see Appendix A), and iterative algorithms based on successive reduction of the mean-squared error [2]-[5]. However, it is well known that this problem is ill-posed, i.e., its solution $f(t)$ is not continuously dependent on $g(t)$, and as Youla [6] says, there exist finite energy perturbations that yield unbounded error. As a result, all the above methods work well if $g(t)$ is exactly a segment of $f(t)$ and there are no errors associated with their numerical implementation. Unfortunately, these requirements are almost never satisfied. For the solutions based on the prolate spheroidal functions expansion, there are available various techniques for stabilizing them. All of them make use of additional knowledge about $f(t)$ beyond its band-limited character. These constraints are usually referred to as "regularizers" of ill-posed problems. In order to obtain such a regularized solution, Slepian [1] assumed that the energy of the unknown signal $f(t)$ is given. Viano [7] derived another approximating solution of $f(t)$ by simultaneously imposing bounds on the energies of the signal $f(t)$ and the data-error. Miller [8] has also considered various energy bound regularizers in an abstract Hilbert space formulation. However, all the above solutions involved the computation and storage of the prolate spheroidal functions, an extremely difficult numerical task. In addition, they are also limited by series truncation

[^0]errors. Similar results have been presented by Bertero et al. [9] and Rushford et al. [10]. On the iterative algorithm, Papoulis [4] suggested the early termination of the iteration but the determination of a terminating criterion is not available yet. Youla [6] investigates this iterative algorithm and provides us with the necessary and sufficient conditions for its convergence.

In this work we modify the iterative algorithm [2], [3] to include energy constraints, connecting in this way the regularized solutions obtained in terms of the prolate spheroidal functions, with the iterative algorithm. The suggested iteration is stable even in the presence of noise and converges into a function which is optimum in minimizing some functionals stated in Section II. Three extrapolating problems with different energy constraints are stated in the next section, and their solutions in terms of the prolate spheroidal functions are presented. In Section III we provide a simple iterative algorithm to obtain the solutions in Section II by using only Fourier transformations. This iteration was obtained by a proper modification of the algorithm proposed by Papoulis and Gerchberg, and it invokes a regularization parameter $\mu$. Section IV discusses the numerical evaluation of $\mu$.

Section $V$ is devoted to the solution of the discrete versions of the three extrapolation problems and an equivalent discrete iteration is also derived. The connection between the continuous extrapolation algorithm and its numerical implementation is also discussed in Section V. Finally, the applicability of the suggested iteration is illustrated with a numerical example presented in Section VI.

## II. Extrapolation with Energy Constraints

To avoid lengthy formulas we adopt the usual norm notation, denoting hereafter

$$
\begin{aligned}
& \|f(t)\|^{2}=\int_{-\infty}^{\infty}|f(t)|^{2} d t \\
& \|f(t)\|_{T}^{2}=\int_{-T}^{T}|f(t)|^{2} d t .
\end{aligned}
$$

For simplicity, throughout this work time functions shall be assumed to be real, but the results are valid also for complex functions.
Let $f(t)$ be a $\sigma$-band-limited function, i.e.,

$$
\begin{align*}
& F(\omega)=0 \quad \text { for } \quad|\omega|>\sigma \\
& \|f(t)\|^{2}<\infty \tag{1}
\end{align*}
$$

where $F(\omega)$ is the Fourier transform of $f(t)$ and let

$$
g(t)=f(t)+n(t) \quad|t|<T
$$

We want to estimate $f(t)$ in terms of its noisy segment $g(t)$. As we mentioned in the Introduction, there exist errors $n(t)$ with finite but arbitrary small energy that yield arbitrary large errors on the solution $f(t)$ at any point outside $(-T, T)$. However, if additional information is given on $f(t)$ beyond its band-limited property, the estimation of $f(t)$ can be improved and the ill-posed character of the problem be removed. In this work, we will assume that energy bounds on $f(t)$ and/or $n(t)$ are known. Based upon the given additional information, we state three different versions of the extrapolation problem; and the optimum estimator $f_{*}(t)$ of $f(t)$ is derived in terms of a prolate spheroidal series expansion.
Problem 1: Given that the energy of $f(t)$ is equal to or less than $R^{2}$, i.e.,

$$
\begin{equation*}
\|f(t)\|^{2} \leqslant R^{2} \tag{2}
\end{equation*}
$$

approximate $f(t)$ with $f_{*}(t)$ in $F_{1}$ such that

$$
\begin{equation*}
I_{1}=\left\|f_{*}(t)-g(t)\right\|_{T}^{2}=\min _{F_{1}}\|f(t)-g(t)\|_{T}^{2} \tag{3}
\end{equation*}
$$

where $F_{1}$ is the family of $\sigma$-band-limited functions $f(t)$ satisfying (2).
This is the same problem considered and solved by Slepian et al. [1]. His solution is the same with the one given in Theorem 1.
Problem 2: Given that the energy of the noise in $(-T, T)$ is equal to or less than $\epsilon^{2}$, i.e., $\|n(t)\|_{T}^{2} \leqslant \epsilon^{2}$, find an $f_{*}(t)$ in $F_{2}$ such that

$$
\begin{align*}
& \left\|f_{*}(t)-g(t)\right\|_{T}^{2} \leqslant \epsilon^{2}  \tag{4}\\
& I_{2}=\left\|f_{*}(t)\right\|^{2}=\min _{F_{2}}\|f(t)\|^{2} \tag{5}
\end{align*}
$$

where $F_{2}$ is the family of $\sigma$-band-limited functions $f(t)$ satisfying $\|f(t)-g(t)\|_{T}^{2} \leqslant \epsilon^{2}$.
Problem 3: Given that the energy of $f(t)$ is equal to or less than $R^{2}$ and that the energy of $n(t)$ in $(-T, T)$ is equal to or less than $\epsilon^{2}$, find a $\sigma$-band-limited function $f_{*}(t)$ such that

$$
\begin{align*}
& \left\|f_{*}(t)-g(t)\right\|_{T}^{2} \leqslant \epsilon^{2}  \tag{6}\\
& \left\|f_{*}(t)\right\|^{2} \leqslant R^{2}
\end{align*}
$$

Problem 3 has been considered also by Viano [7], but he obtained an approximating solution. Finally, Miller [8] has also considered all the above three cases in an abstract Hilbert space and derived similar solutions using the method of eigenfunction expansion [1], but the form of his solutions is different from ours.
Theorem 1: Let

$$
\begin{equation*}
g(t)=\sum_{k=0}^{\infty} b_{k} \phi_{k}(t) \quad \text { for } \quad|t|<T \tag{8}
\end{equation*}
$$

where $\phi_{k}(t)$ are the prolate spheroidal functions (see A1) and $b_{k}$ defined as in (A8). Then a solution to Problems 1-3 is of the form

$$
\begin{equation*}
f_{*}(t)=\sum_{k=0}^{\infty} \frac{\lambda_{k} b_{k}}{\lambda_{k}+\mu} \phi_{k}(t) \tag{9}
\end{equation*}
$$

a) For Problem $1, \mu=0$ or $\mu=\mu_{R}>0$ where $\mu_{R}$ is the solution of

$$
\begin{equation*}
R\left(\mu_{R}\right)=\sum_{k=0}^{\infty} \frac{\lambda_{k}^{2} b_{k}^{2}}{\left(\lambda_{k}+\mu_{R}\right)^{2}}=R^{2} \tag{10}
\end{equation*}
$$

b) For Problem $2, \mu=\infty$ or $0 \leqslant \mu=\mu_{\epsilon}<\infty$ where $\mu_{\epsilon}$ is the solution of

$$
\begin{equation*}
J\left(\mu_{\epsilon}\right)=\sum_{k=0}^{\infty} \frac{\lambda_{k} \mu_{\epsilon}^{2} b_{k}^{2}}{\left(\lambda_{k}+\mu_{\epsilon}\right)^{2}}=\epsilon^{2} \tag{11}
\end{equation*}
$$

c) For Problem 3,

$$
\begin{equation*}
0 \leqslant \mu_{R} \leqslant \mu \leqslant \mu_{\epsilon} \tag{12}
\end{equation*}
$$

If $\mu_{\epsilon}<\mu_{R}$, then Problem 3 does not have a solution: $\mu_{\epsilon}$ and $\mu_{R}$ are the solutions of $R\left(\mu_{R}\right)=R^{2}$ and $J\left(\mu_{\epsilon}\right)=\epsilon^{2}$, respectively.

Proof:
a) With $f(t)$ a $\sigma$-band-limited function (1), we obtain from (A6)

$$
\begin{equation*}
f(t)=\sum_{k=0}^{\infty} a_{k} \phi_{k}(t) \tag{13}
\end{equation*}
$$

Using relations (8), (13), and (A9), we get

$$
\begin{equation*}
I=\|f(t)-g(t)\|_{T}^{2}=\sum_{k=0}^{\infty} \lambda_{k}\left(b_{k}-a_{k}\right)^{2} \tag{14}
\end{equation*}
$$

Thus, Problem 1 is equivalent to determine $a_{k}$, minimizing $I$ under the constraint

$$
\begin{equation*}
\sum_{k=0}^{\infty} a_{k}^{2} \leqslant R^{2} \tag{15}
\end{equation*}
$$

Using Lagrangian multipliers we have the unconstrained minimization of

$$
\begin{equation*}
H\left(a_{k}, \mu\right)=\sum_{k=0}^{\infty} \lambda_{k}\left(b_{k}-a_{k}\right)^{2}+\mu \sum_{k=0}^{\infty} a_{k}^{2} . \tag{16}
\end{equation*}
$$

Setting

$$
\begin{equation*}
\frac{\partial H\left(a_{k}, \mu\right)}{\partial a_{k}}=0 \tag{17}
\end{equation*}
$$

we obtain

$$
\begin{equation*}
a_{k}=\frac{\lambda_{k} b_{k}}{\lambda_{k}+\mu} \tag{18}
\end{equation*}
$$

Substituting (18) into (14) and (15), we have to determine $\mu$ such that

$$
\begin{equation*}
R(\mu)=\sum_{k=0}^{\infty} \frac{\lambda_{k}^{2} b_{k}^{2}}{\left(\lambda_{k}+\mu\right)^{2}} \leqslant R^{2} \tag{19}
\end{equation*}
$$

and minimizing

$$
\begin{equation*}
J(\mu)=\sum_{k=0}^{\infty} \frac{\lambda_{k} \mu^{2} b_{k}^{2}}{\left(\lambda_{k}+\mu\right)^{2}} . \tag{20}
\end{equation*}
$$



Fig. 1. Sketch of $R(\mu)$ and $J(\mu)$.
It can be shown that since for any positive $\mu, J(-\mu)>J(\mu)$ and $R(-\mu)>R(\mu)$, we need to consider only $\mu \geqslant 0$. Since for $\mu \geqslant 0, R(\mu)$ and $J(\mu)$ (Section IV) are, respectively, monotonically decreasing and increasing functions of $\mu$ (see Fig. 1), we have

$$
\mu= \begin{cases}\mu_{R}>0 & \text { if } R^{2}<\sum_{k=0}^{\infty} b_{k}^{2}  \tag{21}\\ 0 & \text { otherwise }\end{cases}
$$

where $\mu_{R}$ is the positive root of $R(\mu)=R^{2}$.
b) Following a procedure which is similar to a), we obtain that the $\mu$ for Problem 2 is

$$
\mu= \begin{cases}\mu_{\epsilon} \geqslant 0 & \text { if } \epsilon^{2}<\|g(t)\|_{T}^{2}  \tag{22}\\ \infty & \text { otherwise }\end{cases}
$$

where $\mu_{\epsilon}$ is the positive root of $J(\mu)=\epsilon^{2}$.
c) If $f_{*}(t)$ is a solution of Problem 3, then

$$
\begin{align*}
J\left(\mu_{R}\right)= & \min _{F_{1}}\|f(t)-g(t)\|_{T}^{2} \leqslant\left\|f_{*}(t)-g(t)\right\|_{T}^{2} \\
& \leqslant \epsilon^{2}=J\left(\mu_{\epsilon}\right) \tag{23}
\end{align*}
$$

because $f_{*}(t) \epsilon F_{1}$. Therefore, since $J(\mu)$ is monotonically increasing, the necessary condition for the existence of a solution to Problem 3 is

$$
\mu_{\epsilon} \geqslant \mu_{R}
$$

It is easy to find that any $\mu$ in the closed interval $\left[\mu_{R}, \mu_{\epsilon}\right]$ will provide us with an acceptable $f_{*}(t)$.

The form obtained in (9) is usually referred to as a "regularized" solution [12] of the original ill-posed extrapolation problem because it is defined for any given segment $g(t)$ and in the absence of noise, i.e., $\mu=0, f_{*}(t)=f(t) . \mu$ is called the regularization parameter.

Also, we would like to comment on the meaning of $\mu=\infty$ for Problem 2. If $\|g(t)\|_{T}^{2} \leqslant \epsilon^{2}$, then it is clear that the trivial solution $f_{*}(t)=0$ is the optimum one because it is $\sigma$-bandlimited, satisfies condition 4 , and possesses minimum energy in $F_{2}$.

The numerical evaluation of the derived optimum solution $f_{*}(t)$ presents two difficulties. The first one is the use of the prolate spheroidal functions and the other one is the calculation of the parameter $\mu$ from (21) and (22). These two problems are addressed in Sections III and IV.

## III. The Iterative Solution

To avoid the complexity associated with the eigenfunctions expansion, i.e., storing $\phi_{k}(t)$, evaluation of $\lambda_{k}$, etc., we shall devise an iteration to converge into $f_{*}(t)$ [see (9)], assuming that $\mu$ is given.
Let

$$
\begin{align*}
f_{n+1}(t)= & {\left[(1-\alpha \mu) f_{n}(t)+\alpha\left(g(t)-f_{n}(t)\right) P_{T}(t)\right] } \\
& * \frac{\sin \sigma t}{\pi t} \tag{24}
\end{align*}
$$

$$
f_{0}(t)=0
$$

where

$$
P_{T}(t)= \begin{cases}1 & |t|<T \\ 0 & |t|>T\end{cases}
$$

and $*$ denotes convolution.
Then, with

$$
0<\alpha<\frac{2}{1+\mu}
$$

$$
\lim _{n \rightarrow \infty} f_{n}(t)=f_{*}(t)
$$

Proof: Because $f_{n}(t)$ is $\sigma$-band-limited,

$$
\begin{equation*}
f_{n}(t)=\sum_{k=0}^{\infty} a_{k, n} \phi_{k}(t) \tag{25}
\end{equation*}
$$

Substituting $f_{n}(t)$ and $f_{n+1}(t)$ from (25) and $g(t)$ from (8) into (24), and using (A1) and (A2), we readily obtain, by equating the coefficients of $\phi_{k}(t)$,

$$
\begin{equation*}
a_{k, n+1}=\left[1-\alpha\left(\mu+\lambda_{k}\right)\right] a_{k, n}+\alpha \lambda_{k} b_{k} \tag{26}
\end{equation*}
$$

which is a recursive equation in $a_{k, n}$ with $a_{k, 0}=0$. Its solution is

$$
\begin{equation*}
a_{k, n}=\frac{\lambda_{k} b_{k}}{\lambda_{k}+\mu}\left[1-\left(1-\alpha\left(\lambda_{k}+\mu\right)\right)^{n}\right] \tag{27}
\end{equation*}
$$

Setting

$$
\begin{equation*}
0<\alpha<\frac{2}{1+\mu}<\frac{2}{\lambda_{k}+\mu} \tag{28}
\end{equation*}
$$

we obtain [see (A7) and (27)]

$$
\begin{align*}
\lim _{n \rightarrow \infty} & \left\|f_{n}(t)-f_{*}(t)\right\|^{2} \\
& =\lim _{n \rightarrow \infty} \sum_{k=0}^{\infty}\left(\frac{\lambda_{k} b_{k}}{\lambda_{k}+\mu}\right)^{2}\left(1-\alpha\left(\lambda_{k}+\mu\right)\right)^{2 n} \\
& =0 \tag{29}
\end{align*}
$$

which results in

$$
\begin{equation*}
\lim _{n \rightarrow \infty} f_{n}(t)=f_{*}(t) \tag{30}
\end{equation*}
$$

because for the $\sigma$-band-limited function, $f_{n}(t)-f_{*}(t)[11, \mathrm{p}$. 169],

$$
\max \left|f_{n}(t)-f_{*}(t)\right| \leqslant \sqrt{\frac{\sigma}{\pi}}\left\|f_{n}(t)-f_{*}(t)\right\|
$$

This completes the proof because of (29).

## Notes

1) The constant $\alpha$ is usually referred to as a relaxation parameter. It is used to guarantee the convergence of the iteration. If monotonic convergence, i.e., $\left\|f_{n}(t)\right\|>\left\|f_{n+1}(t)\right\|$, is desirable, then the condition $0<\alpha \leqslant 1 / 1+\mu$ is sufficient. This can be easily obtained from (27) because

$$
\begin{aligned}
\left\|f_{n}(t)\right\|^{2} & =\sum_{k=0}^{\infty}\left|a_{k, n}\right|^{2}>\sum_{k=0}^{\infty}\left|a_{k, n-1}\right|^{2} \\
& =\left\|f_{n-1}(t)\right\|^{2}
\end{aligned}
$$

2) For $\mu=0$ and $\alpha=1$ the iteration in (24) becomes the Papoulis-Gerchberg algorithm. This iteration converges if

$$
\sum_{k=0}^{\infty} b_{k}^{2}<\infty
$$

which is equivalent to saying that the noise $n(t)$ in (1) is a segment of a $\sigma$-band-limited function.
The significance of the parameter $\mu$ in the iteration can also be described in terms of integral equations.

If we rearrange the terms in (24), and make use of the fact that $f_{n}(t)$ is $\sigma$-band-limited, we obtain

$$
\begin{align*}
f_{n+1}(t)= & (1-\alpha \mu) f_{n}(t)+\alpha \int_{-T}^{T}(g(\tau) \\
& \left.-f_{n}(\tau)\right) \frac{\sin \sigma(t-\tau)}{\pi(t-\tau)} d \tau \tag{31}
\end{align*}
$$

Because $f_{n}(t) \rightarrow f_{*}(t)$ uniformly (30), taking limits for $n \rightarrow \infty$ on both sides of (31), we have that $f_{*}(t)$ should satisfy

$$
\begin{align*}
\mu f_{*}(t) & +\int_{-T}^{T} f_{*}(\tau) \frac{\sin \sigma(t-\tau)}{\pi(t-\tau)} d \tau \\
= & \int_{-T}^{T} g(\tau) \frac{\sin \sigma(t-\tau)}{\pi(t-\tau)} d \tau \tag{32}
\end{align*}
$$

But this is a Fredholm integral equation of the second kind, and it is well known [12] that the solution $f_{*}(t)$ of this equation is well behaved. It can be shown also directly from (9) that $f_{s}(t)$ is the solution of (32) and then the algorithm in (24) can be simply considered as the successive approximation solution of (32).
Depending on the numerical implementation of the iterative solution, (31) may be a more preferable form than (24), because it avoids the use of infinite integrals. If $\alpha=1$ and $\mu=$

0 , (31) becomes the iteration used by Cadzow in [13]. Notice that in this case the solution $f_{*}(t)$ satisfies a Fredholm integral equation of the first kind.

$$
\begin{equation*}
\int_{-\infty}^{\infty} f_{*}(\tau) \frac{\sin \sigma(t-\tau)}{(t-\tau)} d \tau=f(t) \quad t \in[-T, T] \tag{33}
\end{equation*}
$$

The following corollary shows the stability of the iterative solution (24).
Corollary 1: For any $\mu>0$ and $0<\alpha \leqslant \frac{2}{1+\mu}$

$$
\begin{equation*}
\left\|f_{n}(t)\right\|^{2}<\frac{\|g(t)\|_{T}^{2}}{\mu^{2}} \tag{34}
\end{equation*}
$$

Proof: From (27) we conclude that

$$
\left\|f_{n}(t)\right\|^{2}<\left\|f_{*}(t)\right\|^{2}
$$

and (34) results from the fact that

$$
\begin{aligned}
\left\|f_{*}(t)\right\|^{2}= & \sum_{k=0}^{\infty} \frac{\lambda_{k}^{2} b_{k}^{2}}{\left(\lambda_{k}+\mu\right)^{2}}<\sum_{k=0}^{\infty} \frac{\lambda_{k}^{2} b_{k}^{2}}{\mu^{2}} \\
& <\frac{1}{\mu^{2}} \sum_{k=0}^{\infty} \lambda_{k} b_{k}^{2}=\frac{1}{\mu^{2}}\|g(t)\|_{T}^{2}
\end{aligned}
$$

where for the inequalities we used that $0<\lambda_{k}<1$ and for the equality (A9). Actually, we can also prove that

$$
\begin{equation*}
\left\|f_{n}(t)\right\|^{2}<\frac{\|g(t)\|_{T}^{2}}{2 \mu} \tag{35}
\end{equation*}
$$

## IV. On the Evaluation of $\mu_{R}$ and $\mu_{\epsilon}$

The evaluation of $\mu_{R}$ and $\mu_{\epsilon}$ is the second problem we have to consider. Relations (10) or (11) cannot be used because they require knowledge of $\lambda_{k}$ and $b_{k}$. To avoid this problem we use the iteration to obtain $R(\mu)$ or $J(\mu)$ for various values of $\mu$ and then utilize one of the various searching numerical techniques (for example, bisection or secant, etc.) to find the solution of (11) or (12). Below we present some theorems which can be used to speed up the evaluation of $\mu_{R}$ or $\mu_{\epsilon}$.
Lemma 1: Let $f_{*}(t, \mu)$ denote the limit of the iteration (24) with a given $\mu$, then

$$
\begin{equation*}
\left\|f_{*}(t, \mu)\right\|_{T}^{2}+2 \mu R(\mu)+J(\mu)=\|g(t)\|_{T}^{2} \tag{36}
\end{equation*}
$$

Proof: Using relations (19), (20), and

$$
\left\|f_{*}(t ; \mu)\right\|_{T}^{2}=\sum_{k=0}^{\infty} \frac{\lambda_{k}^{3} b_{k}^{2}}{\left(\lambda_{k}+\mu\right)^{2}}
$$

we obtain

$$
\begin{align*}
& \left\|f_{*}(t, \mu)\right\|_{T}^{2}+2 \mu R(\mu)+J(\mu) \\
& =\sum_{k=0}^{\infty}\left(\lambda_{k}^{3}+2 \mu \lambda_{k}^{2}+\lambda_{k} \mu^{2}\right) \frac{b_{k}^{2}}{\left(\lambda_{k}+\mu\right)^{2}} \\
& =\sum_{k=0}^{\infty} \lambda_{k} b_{k}^{2}=\|g(t)\|_{T}^{2} .
\end{align*}
$$

Lemma 2: The functions $R(\mu)$ and $J(\mu)$ for $\mu \geqslant 0$ are monotonically decreasing and increasing, respectively.

Proof: Taking the derivatives we can easily show (19), (20) that for $\mu>0$,

$$
\begin{aligned}
R^{\prime}(\mu) & =-2 \sum_{k=0}^{\infty} \frac{\lambda_{k}^{2} b_{k}^{2}}{\left(\lambda_{k}+\mu\right)^{3}}<0
\end{aligned} \quad \mu>00
$$

Theorem 3: The constants $\mu_{R}$ and $\mu_{\epsilon}$ satisfy the following bounds:

$$
\begin{align*}
& 0 \leqslant \mu_{R}<\frac{\|g(t)\|_{T}^{2}}{2 R^{2}}  \tag{37}\\
& 0 \leqslant \mu_{\epsilon}<\frac{\epsilon}{\|g(t)\|_{T}-\epsilon} \quad \text { if } \quad\|g(t)\|_{T}^{2}>\epsilon^{2} \tag{38}
\end{align*}
$$

Proof: The positivity of $\mu_{R}$ and $\mu_{\epsilon}$ has been proved in Theorem 1. The upper bound of (37) is obtained directly from (36) because $R\left(\mu_{R}\right)=R^{2}$. From (36) we also obtain, assuming $\mu<\infty$,

$$
\frac{\|g(t)\|_{T}^{2}-J(\mu)}{1+2 \mu}<R(\mu)<\sum_{k=0}^{\infty} \frac{\lambda_{k} b_{k}^{2}}{\left(\lambda_{k}+\mu\right)^{2}}=\frac{J(\mu)}{\mu^{2}}
$$

or

$$
\mu<\frac{\sqrt{J(\mu)}}{\|g(t)\|_{T}-\sqrt{J(\mu)}}
$$

and (38) follows because $J\left(\mu_{\epsilon}\right)=\epsilon^{2}$.
Theorem 3 provides us with an initial value $\mu_{0}$ for $\mu_{\epsilon}$ or $\mu_{R}$. Using the iteration we can evaluate $R\left(\mu_{0}\right)=\left\|f_{*}\left(t, \mu_{0}\right)\right\|^{2}$ or $J\left(\mu_{0}\right)=\left\|g(t)-f_{*}\left(t, \mu_{0}\right)\right\|_{T}^{2}$; and utilizing the monotonicity of $R(\mu)$ and $J(\mu)$, we can evaluate the next value $\mu_{1}$ by using various numerical methods for solving nonlinear equations.

$$
\begin{align*}
& \text { Theorem 4: With } 0<\alpha \leqslant \frac{1}{1+\mu}  \tag{39}\\
& R(\mu)\left[1-\left(\frac{1}{1+\mu}\right)^{n}\right]^{2}<\left\|f_{n}(t)\right\|^{2}<R(\mu)  \tag{40}\\
& J(\mu) \leqslant\left\|f_{n}(t)-g(t)\right\|_{T}^{2}<J(\mu)\left[1+\frac{1}{\mu}(1-\alpha \mu)^{n}\right]^{2} . \tag{41}
\end{align*}
$$

Proof: From (27) we obtain

$$
\begin{equation*}
\left.\left\|f_{n}(t)\right\|^{2}=\sum_{k=0}^{\infty} \frac{\lambda_{k}^{2} b_{k}^{2}}{\left(\lambda_{k}+\mu\right)^{2}}\left[1-\alpha\left(\lambda_{k}+\mu\right)\right)^{n}\right]^{2} \tag{42}
\end{equation*}
$$

Using (39) and $0<\lambda_{k}<1$ we readily obtain

$$
\begin{aligned}
1 & >\left(1-\left(1-\alpha\left(\lambda_{k}+\mu\right)\right)^{n}\right)^{2}>\left(1-\left(1-\frac{\lambda_{k}+\mu}{1+\mu}\right)^{n}\right)^{2} \\
& =\left[1-\left(\frac{1-\lambda_{k}}{1+\mu}\right)^{n}\right]^{2}>\left[1-\left(\frac{1}{1+\mu}\right)^{n}\right]^{2}
\end{aligned}
$$

and (40) follows from the above and (19). Similarly, we can prove (41) because

$$
\begin{align*}
& \left\|f_{n}(t)-g(t)\right\|_{T}^{2} \\
& \quad=\sum_{k=0}^{\infty} \frac{\lambda_{k} b_{k}^{2} \mu^{2}}{\left(\lambda_{k}+\mu\right)^{2}}\left(1+\frac{\lambda_{k}}{\mu}\left(1-\alpha\left(\lambda_{k}+\mu\right)\right)^{n}\right)^{2} \tag{43}
\end{align*}
$$

Relation (40) can reduce significantly the number of calculations in Problem 1, because

$$
\text { if }\left\|f_{n}(t)\right\|^{2}>R^{2}, \quad \text { then } \quad \mu>\mu_{R}
$$

and

$$
\begin{equation*}
\text { if }\left\|f_{n}(t)\right\|^{2} R^{2}\left(1-\left(\frac{1}{1+\mu}\right)^{n}\right)^{2}, \quad \text { then } \quad \mu<\mu_{R} \tag{44}
\end{equation*}
$$

Similarly, using (41) we have for Problem 2

$$
\text { if }\left\|f_{n}(t)-g(t)\right\|_{T}^{2}<\epsilon^{2}, \quad \text { then } \quad \mu<\mu_{\epsilon}
$$

and

$$
\begin{align*}
& \text { if }\left\|f_{n}(t)-g(t)\right\|_{T}^{2}>\epsilon^{2}\left(1+\frac{1}{\mu}\left(1-\alpha\left(\lambda_{k}+\mu\right)\right)^{n}\right)^{2} \\
& \text { then } \mu>\mu_{\epsilon} \tag{45}
\end{align*}
$$

Finally; (40) and (41) can be used simultaneously to help us determine an acceptable value of $\mu$ for Problem 3.
The application of the above statements will be demonstrated in Section VI with a numerical example.

## V. The Discrete Problem

There recently has developed a controversy over whether analog algorithms can be digitized, and over the connection between the results of a numerical implementation versus the analog solution. For discrete signals the analyticity property vanishes due to sampling, and extrapolated estimate need not coincide with the original one. However, if the energy bounds, described in Section II, are imposed, the uniqueness of the solution is restored and the connection between the analog solution and its discrete digital implementation can be established.
In this section we will restate the problems of Section II in a discrete environment and rederive their solutions. In Section V-A we will consider the problem of extrapolating a bandlimited sequence $f[n]$ for any $n$, in terms of a finite set of noisy samples. This problem is obtained when a $\sigma$-band-limited function $f(t)$ is sampled with a sampling interval $T_{s} \leqslant \pi / \sigma$, and we want to estimate its Fourier transform in terms of a finite set of samples $f[n]=f\left(n T_{s}\right)|n| \leqslant M$. In Section V-B the problem of extrapolating $N$-periodic band-limited sequences will be presented and its solution under the equivalent energy constraints will be given. This situation arises every time the environment of a computer is used. Digitization and roundoff errors can be considered as part of our noise term. Finally, in Section V-C the relation between the analog algorithm and its numerical realization is derived.

Also, we would like to mention that in addition to the importance of this section to the numerical implementation
of the analog algorithm, there are cases where the problems are phrased directly in their discrete form (cases A and B), i.e., we are asked to extrapolate periodic functions, or trigonometric polynomials in terms of their samples.
Prolate spheroidal functions (PSF) were the basic tool of the analog solution in Section II. Similarly, the key function for the two discrete problems will be the discrete prolate spheroidal sequences (DPSS) and the periodic discrete prolate spheroidal sequence (P-DPSS). Their basic properties and definitions are stated in Appendixes B and C. For a more detailed description of DPSS see Slepian [14] or Papoulis [15], and for P-DPSS see [16] and [17].

Throughout this work the usual norm notation for sequences is used, denoting hereafter

$$
\begin{aligned}
& \|f[n]\|^{2}=\sum_{n=-\infty}^{\infty}|f[n]|^{2} \\
& \|f[n]\|_{M, N}^{2}=\sum_{n=M}^{N}|f[n]|^{2} .
\end{aligned}
$$

## A. Extrapolating Band-Limited Sequences with Energy Constraints

Let $f[n]$ be a $\sigma$-band-limited sequence, i.e.,

$$
\begin{aligned}
& F(\theta)=\sum_{n=-\infty}^{\infty} f[n] e^{-j n \theta}=0 \quad \pi>|\theta|>\sigma \\
& \|f[n]\|^{2}<\infty
\end{aligned}
$$

where $F(\theta)$ is the discrete Fourier transform of $f[n]$ and let

$$
g[n]=f[n]+\eta[n] \quad|n| \leqslant M
$$

be the given data. We want to estimate $f[n]$, for all $n$, in terms of its noisy segment $g[n]$.

It should be pointed out again that in constrast with the continuous case, $f[n]$ cannot be determined uniquely from $g[n]$ in the absence of noise. However, imposing the discrete version of the constraints (2), (3), or (4) and (5), the uniqueness of the solution can be reestablished, and a theorem equivalent to Theorem 1 will be derived.
Lemma 3: Let $h[m]$ be known for $|m| \leqslant M$. Then the function $f^{+}[m]$ satisfying the conditions

$$
\begin{aligned}
& f^{+}[m] \text { is } \sigma \text {-band-limited } \\
& f^{+}[m]=h[m] \quad|m| \leqslant M
\end{aligned}
$$

and $\left\|f^{+}[m]\right\|^{2}=\min \left\{\|f[m]\|^{2} ; f[m]\right.$ is $\sigma$-band-limited and $f[m]=h[m]|m| \leqslant M\}$ is given by

$$
\begin{equation*}
f^{+}[m]=\sum_{k=0}^{2 M} a_{k} \phi_{k}[m] \tag{46}
\end{equation*}
$$

where $\phi_{k}[m]$ are the DPSS described in Appendix B; and

$$
a_{k}=\frac{1}{\lambda_{k}} \sum_{m=-M}^{M} h[m] \phi_{k}[m] .
$$

Proof: Let $f[m]$ be a $\sigma$-band-limited sequence and $f[m]=$ $h[m]$ for $|m| \leqslant M$. Then, from (B6)

$$
\begin{align*}
f[m] & =\sum_{k=0}^{\infty} a_{k} \phi_{k}[m] \\
& =\sum_{k=0}^{2 M} a_{k} \phi_{k}[m]+\sum_{k=2 M+1}^{\infty} a_{k} \phi_{k}[m] \\
& =f^{+}[m]+f_{1}[m] \tag{47}
\end{align*}
$$

where $\phi_{k}[m], k=0,1, \cdots, 2 M$ are the DPSSS, but

$$
f_{1}[m]=0 \quad \text { for } \quad|m| \leqslant M
$$

[see (B8)], thus,

$$
\begin{equation*}
f[m]=f^{+}[m] \quad|m| \leqslant M \tag{48}
\end{equation*}
$$

and

$$
\left\|f^{+}[m]\right\|^{2} \leqslant\|f[m]\|^{2}
$$

This result is used also in [17] where $f^{+}[m]$ is referred to as the "minimum norm least square solution."

Definition: The linear $(2 M+1)$-dimensional subspace determined by $\phi_{k}[m], k=0,1, \cdots, 2 M$, where $\phi_{k}[m]$ are the $(2 M+1)$ DPSS, is denoted by $F_{3}$.

Using Lemma 3, the space of $\sigma$-band-limited sequences in Problems 1-3 can be replaced by $F_{3}$. Then the uniqueness of the solution has been restored and all the proofs in Sections IIIV can easily be duplicated for the $\sigma$-band-limited sequences case; and the equivalent results are obtained by merely substituting $\infty$ by $2 M$ and $f(t)$ by $f[m]$ in Sections II-IV. For example, $f_{*}[m]$, the equivalent of $f_{*}(t)$ in Theorem 1 , is

$$
\begin{equation*}
f_{*}[m]=\sum_{k=0}^{2 M} \frac{\lambda_{k} b_{k}}{\lambda_{k}+\mu} \phi_{k}[m] \tag{49}
\end{equation*}
$$

where $\phi_{k}[m]$ are the DPSS [see (B1), and $b_{k}$ is defined as in (B8)]. At this point, it should be pointed out that the special case of Problem 2 with $\epsilon=0$ has been also considered and solved in [17].
The equivalent discrete form of the iterative solution is easily found to be

$$
\begin{align*}
f_{n+1}[m]= & {\left[(1-\alpha \mu) f_{n}[m]+\alpha\left(g[m]-f_{n}[m]\right) P_{M}[m]\right] } \\
& * \frac{\sin \sigma m}{\pi m} \tag{50}
\end{align*}
$$

with $f_{0}[m]=0$ where $*$ denotes discrete convolution, and

$$
P_{M}[m]= \begin{cases}0 & |m|>M  \tag{51}\\ 1 & |m| \leqslant M\end{cases}
$$

An easier way to implement (49) [see (31)] is

$$
\begin{align*}
f_{n+1}[m]= & (1-\alpha \mu) f_{n}[m] \\
& +\alpha \sum_{l=-M}^{M}\left(g[l]-f_{n}[l]\right) \frac{\sin \sigma(m-l)}{\pi(m-l)} \tag{52}
\end{align*}
$$

because it avoids the infinite summation. Similar to (32), the solution $f_{*}[m]$ in (49) satisfies the equation

$$
\begin{gather*}
\mu f_{*}[m]+\sum_{l=-M}^{M} f_{*}[l] \frac{\sin \sigma(m-l)}{\pi(m-l)} \\
=\sum_{l=-M}^{M} g[l] \frac{\sin \sigma(m-l)}{\pi(m-l)} \tag{53}
\end{gather*}
$$

Besides the iterative algorithms mentioned above [(50) or (52)], (53) can be solved in the following way, which is similar to that provided in [17] and [18] for the noiseless case.
Denote

$$
\delta[m-l]= \begin{cases}1 & m=l \\ 0 & m \neq l\end{cases}
$$

First, solve the $(2 M+1)$ Toeplitz equations

$$
\begin{align*}
& \sum_{l=-M}^{M}\left[\mu \delta[m-l]+\frac{\sin \sigma(m-l)}{\pi(m-l)}\right]\left[g[l]-f_{*}[l]\right] \\
& \quad=\mu g[m] \quad|m| \leqslant M \tag{54}
\end{align*}
$$

and determine $f_{*}[l]$ for $|l| \leqslant M$. Then obtain the solution $f_{*}[m]$ for any $m$ by using

$$
\begin{equation*}
f_{*}[m]=\frac{1}{\mu} \sum_{l=-M}^{M}\left[g[l]-f_{*}[l]\right] \frac{\sin \sigma(m-l)}{\pi(m-l)} \tag{55}
\end{equation*}
$$

Compare to the noiseless case [17], [18], where the corresponding Toeplitz matrix is

$$
\left[\frac{\sin \sigma(m-l)}{\pi(m-l)}\right]_{m, l=-M, \cdots, 0, \cdots, M}
$$

which is an ill-conditioned matrix when $M$ is large. In (54), however, the ill-conditioned character of the matrix is improved because of the addition of the constant $\mu$ to its diagonal elements. This is a common technique for stabilizing an illconditioned matrix; however, here $\mu$ is related to the imposed energy bounds on $f[m]$ and/or $\eta[m]$.

## B. Extrapolating Periodic Band-Limited Sequence with Energy Constraints

Let $f[n]$ be a $K$-band-limited $N$-periodic sequence, i.e., $f[n]=f[n+N]$ and

$$
\begin{align*}
F[m] & =\frac{1}{N} \sum_{n=0}^{N-1} f[n] \exp (-j 2 \pi m n / N) \\
& =0 \quad K<m<N-K \tag{56}
\end{align*}
$$

where $F[m]$ is the discrete Fourier series of the $N$-periodic sequence $f[n]$, and let

$$
g[n]=f[n]+\eta[n] \quad|n| \leqslant M
$$

be the given data, where $2 M+1<N$. We want to estimate the $N$-periodic sequence $f[n]$, in terms of its noisy segment $g[n]$.

Here we should examine the cases of $K \geqslant M$ and $K<M$ separately.
i) For the case $K \geqslant M$, all the statements will be similar to
those in Section V-A. For example, the equivalent solution of Theorem 1 is

$$
\begin{equation*}
f_{*}[m]=\sum_{i=0}^{2 M} \frac{\lambda_{i} b_{i}}{\lambda_{i}+\mu} \phi_{i}[m] \tag{57}
\end{equation*}
$$

where $\phi_{i}[m]$ are the P-DPSS [see (C1)] and $b_{i}$ is defined as in (C12).
The iterative solution is

$$
\begin{align*}
f_{n+1}[m]= & \sum_{l=l_{0}}^{l_{0}+N-1}\left[(1-\alpha \mu) f_{n}[l]+\alpha(g[l]\right. \\
& \left.\left.-f_{n}[l]\right) P_{M}[l]\right] \frac{\sin \frac{\pi}{N}(2 K+1)(m-l)}{N \sin \frac{\pi}{N}(m-l)}
\end{align*}
$$

with

$$
f_{0}[m]=0
$$

or

$$
\begin{array}{r}
f_{n+1}[m]=(1-\alpha \mu) f_{n}[m]+\alpha \sum_{l=-M}^{M}(g[l] \\
\left.-f_{n}[l]\right) \frac{\sin \frac{\pi}{N}(2 K+1)(m-l)}{N \sin \frac{\pi}{N}(m-l)} \tag{59}
\end{array}
$$

with

$$
f_{0}[m]=0
$$

The solution $f_{*}[m]$ satisfies the equation

$$
\begin{gather*}
\mu f_{*}[m]+\sum_{l=-M}^{M} f_{*}[l] \frac{\sin \frac{\pi}{N}(2 K+1)(m-l)}{N \sin \frac{\pi}{N}(m-l)} \\
=\sum_{l=-M}^{M} g[l] \frac{\sin \frac{\pi}{N}(2 K+1)(m-l)}{N \sin \frac{\pi}{N}(m-l)} . \tag{60}
\end{gather*}
$$

ii) For the case $K<M$, see (C14),

$$
\begin{equation*}
g[m]=\sum_{i=0}^{2 K} b_{i} \phi_{i}[m]+\tilde{g}[m] \quad|m| \leqslant M \tag{61}
\end{equation*}
$$

where $\phi_{i}[m]$ are the P-DPSS with nonzero eigenvalues and the term $\tilde{g}=(g[-M], \cdots, g[M])$ is orthogonal to $\phi_{i}=\left(\phi_{i}[-M]\right.$, $\left.\cdots, \phi_{i}[M]\right), i=0, \cdots, 2 K$. For the extrapolating Problem 1 , the solution is

$$
f_{*}[m]=\sum_{i=0}^{2 K} \frac{\lambda_{i} b_{i}}{\lambda_{i}+\mu} \phi_{i}[m]
$$

with $\mu$ the solution of

$$
R(\mu)=\sum_{i=0}^{2 K} \frac{\lambda_{i}^{2} b_{i}^{2}}{\left(\lambda_{i}+\mu\right)^{2}}=R^{2}
$$

For Problem 2, there is no solution if $\|\widetilde{g}[m]\|_{-M, M}^{2}>\epsilon^{2}$. When $\|\tilde{g}[m]\|_{-M, M}^{2} \leqslant \epsilon^{2}$, the solution is

$$
\begin{equation*}
f_{*}[m]=\sum_{i=0}^{2 K} \frac{\lambda_{i} b_{i}}{\lambda_{i}+\mu} \phi_{i}[m] \tag{62}
\end{equation*}
$$

where $\mu=\infty$ if $\epsilon^{2}=\|\tilde{g}[m]\|_{-M, M}^{2}$ or $\mu$ satisfies the equation

$$
\begin{gather*}
J(\mu) \triangleq \sum_{i=0}^{2 K} \frac{\lambda_{i} \mu^{2} b_{i}^{2}}{\left(\lambda_{i}+\mu\right)^{2}}=\epsilon^{2}-\|\tilde{g}[m]\|_{-M, M}^{2} \\
\text { for } \epsilon^{2}>\|\tilde{g}[m]\|_{-M, M}^{2} . \tag{63}
\end{gather*}
$$

Denote the solution of (63) by $\mu_{\epsilon^{\prime}}$ where

$$
\epsilon^{\prime}=\sqrt{\epsilon^{2}-\|\tilde{g}[m]\|_{-M, M}^{2}}
$$

Substituting $\mu_{\epsilon}$ by $\mu_{\epsilon}$, the assertion which is similar to c) in Theorem 1 can be obtained.
When $K<M$, (58) and (59) as well as (60) can also be derived assuming that the solution exists.

## C. Digital Implementation of the Analog Algorithm

Now we will discuss the connection between the analog solution of the algorithm and the solution of its digital implementation. Cadzow [13] has also studied this problem, but only for the noiseless case.
Let $\mu>0$ and $g(t)$ given for $|t| \leqslant T$. If we sample $g(t)$ with a sampling interval $T_{s} \leqslant \pi / \sigma$, we obtain the discrete data

$$
\begin{equation*}
g[l]=g\left(l T_{s}\right) \quad l=-M, \cdots, M, \quad M=T / T_{s} . \tag{64}
\end{equation*}
$$

Let us choose $N$ large enough such that if

$$
K=\text { integer }\left[\begin{array}{ll}
\frac{(N-1)}{2} & \frac{\sigma T_{s}}{\pi}
\end{array}\right]
$$

then $N>2 M+1$ and $K>M$. Applying the iterations described in (58) or (59) on the given data $g$ [l], an $N$-periodic sequence $f_{*}^{N}, T_{s}\left(m T_{s}\right)$ is obtained satisfying the equation

$$
\begin{align*}
& \mu f_{*}^{N, T_{s}\left(m T_{s}\right)+\sum_{l=-M}^{M} f_{*}^{N, T_{s}\left(l T_{s}\right)} \frac{\sin \frac{\pi}{N}(2 K+1)(m-l)}{N \sin \frac{\pi}{N}(m-l)}} \begin{array}{l}
=\sum_{l=-M}^{M} g\left(l T_{s}\right) \frac{\sin \frac{\pi}{N}(2 K+1)(m-l)}{N \sin \frac{\pi}{N}(m-l)} \\
m=-\frac{N}{2}, \cdots, \frac{N}{2}-1 .
\end{array} .
\end{align*}
$$

For simplicity, $N$ is assumed to be even.
For fixed $T_{s}$, let $N$ tend to infinity. Since

$$
\begin{equation*}
\lim _{N \rightarrow \infty} \frac{\sin \frac{\pi}{N}(2 K+1)(m-l)}{N \sin \frac{\pi}{N}(m-l)}=\frac{\sin \sigma T_{s}(m-l)}{\pi(m-l)} \tag{66}
\end{equation*}
$$

if we define

$$
\overline{f_{*}} \bar{T}_{s}\left(m T_{s}\right)= \begin{cases}f_{*}^{N, T_{s}}\left(m T_{s}\right) & m=\frac{N}{2}, \cdots, \frac{N}{2}-1  \tag{67}\\ 0 & \text { otherwise }\end{cases}
$$

then, for sufficiently large $N, \bar{f}_{*}^{N}, T_{s}\left(m T_{s}\right)$ is an approximation of $f_{*}^{T_{s}}\left(m T_{s}\right)$, which satisfies the equation

Now let $T_{s}$ tend to zero. Based on the theory of approximate solution for Fredholm integral equation of the second kind [19], $f_{*}^{T_{s}}\left(m T_{s}\right)$ will be an approximate solution of (32) when $T_{s}$ is sufficiently small.
In short, choosing appropriate $T_{s}$ and $N$ and using iteration (58), the solution of (32) can be derived approximately.

Let us denote respectively with $R^{(i)}(\mu)$ and $J^{(i)}(\mu), i=1,2$, 3 the $R(\mu)$ and $J(\mu)$ of the $\sigma$-band-limited functions, of the $T_{s} \sigma$-band-limited sequences, and of the $\left[(N-1 / 2)\left(\sigma T_{s} / \pi\right)\right]$ -band-limited periodic sequences.

Theorem 5:
i) $\lim _{N \rightarrow \infty} J^{(3)}(\mu)=J^{(2)}(\mu)$
ii) $\lim _{N \rightarrow \infty} R^{(3)}(\mu)=R^{(2)}(\mu)$
iii) $\lim _{T_{s} \rightarrow 0} J^{(2)}(\mu) T_{s}=J^{(1)}(\mu)$
iv) $\lim _{T_{s} \rightarrow 0} R^{(2)}(\mu) T_{s}=R^{(1)}(\mu)$.

Proof: From (65), it follows that
$R^{(3)}(\mu)=\sum_{m=-(N / 2)}^{(N / 2)-1} \mid f_{*}^{N,\left.T_{s}\left(m T_{s}\right)\right|^{2}}$

$$
\begin{aligned}
& =\frac{1}{\mu^{2}} \sum_{m=-(N / 2)}^{(N / 2)-1} \sum_{l=-M}^{M} \sum_{n=-M}^{M} \\
& \cdot\left[g\left(l T_{s}\right)-f_{*}^{\left.N, T_{s}\left(l T_{s}\right)\right]\left[g\left(n T_{s}\right)-f_{*}^{\left.N, T_{s}\left(n T_{s}\right)\right]}\right.}\right. \\
& \quad \cdot \frac{\sin \frac{\pi}{N}(2 K+1)(l-m)}{N \sin \frac{\pi}{N}(l-m)} \frac{\sin \frac{\pi}{N}(2 K+1)(m-n)}{N \sin \frac{\pi}{N}(m-n)}
\end{aligned}
$$

$$
=\frac{1}{\mu^{2}} \sum_{l=-M}^{M} \sum_{n=-M}^{M}\left[g\left(l l_{s}\right)-f_{*}^{\left.N, T_{s}\left(l T_{s}\right)\right]}\right.
$$

$$
\left[g\left(n T_{s}\right)-f_{*}^{\left.N, T_{s}\left(n T_{s}\right)\right]} \frac{\sin \frac{\pi}{N}(2 K+1)(l-n)}{N \sin \frac{\pi}{N}(l-n)}\right.
$$

$$
\begin{align*}
& \mu f_{*}^{T_{s}}\left(m T_{s}\right)+\sum_{l=-M}^{M} f_{*}^{T_{s}}\left(l T_{s}\right) \frac{\sin \sigma T_{s}(m-l)}{T_{s} \pi(m-l)} T_{s} \\
& =\sum_{l=-M}^{M} g\left(l T_{s}\right) \frac{\sin \sigma T_{s}(m-l)}{T_{s} \pi(m-l)} T_{s} \quad|m|=0,1, \cdots . \tag{68}
\end{align*}
$$

Therefore,

$$
\begin{aligned}
\lim _{N \rightarrow \infty} & R^{(3)}(\mu)=\frac{1}{\mu^{2}} \sum_{l=-M}^{M} \sum_{n=-M}^{M}\left[g\left(l T_{s}\right)-f_{*}^{T_{s}}\left(l T_{s}\right)\right] \\
\cdot & {\left[g\left(n T_{s}\right)-f_{*}^{T_{s}}\left(n T_{s}\right)\right] \frac{\sin \sigma T_{s}(l-n)}{\pi(l-n)}=R^{(2)}(\mu) }
\end{aligned}
$$

Furthermore, Lemma 1 still holds for the discrete case. Hence,

$$
\begin{align*}
& J^{(3)}(\mu)-J^{(2)}(\mu)=\left\|f_{*}^{T_{s}}\left(l T_{s}\right)\right\|_{-M, M}^{2} \\
& \quad-\| f_{*}^{N, T_{s}\left(l T_{s}\right) \|_{-M, M}^{2}+2 \mu R^{(2)}(\mu)-2 \mu R^{(3)}(\mu)} \tag{73}
\end{align*}
$$

It is easy to find from (73) and (70) that

$$
\lim _{N \rightarrow \infty} J^{(3)}(\mu)=J^{(2)}(\mu)
$$

Also, using the equation

$$
\begin{gathered}
R^{(2)}(\mu) T_{s}=\frac{1}{\mu^{2}} \sum_{l=-M}^{M} \sum_{n=-M}^{M}\left[g\left(l T_{s}\right)-f_{*}^{\left.T_{s}\left(l T_{s}\right)\right]}\right. \\
\quad \cdot\left[g\left(n T_{s}\right)-f_{*}^{T_{s}}(n s)\right] \frac{\sin \sigma T_{s}(l-n)}{\pi T_{s}(l-n)} T_{s} T_{s}
\end{gathered}
$$

we obtain easily that

$$
\begin{aligned}
& \lim _{T_{s} \rightarrow 0} R^{(2)}(\mu) T_{s}=\frac{1}{\mu^{2}} \int_{-T}^{T} \int_{-T}^{T}\left[g(t)-f_{*}(t)\right] \\
& \quad \cdot\left[g(\tau)-f_{*}(\tau)\right] \frac{\sin \sigma(t-\tau)}{\pi(t-\tau)} d t d \tau=R^{(1)}(\mu)
\end{aligned}
$$

Using Lemma 1 , it follows that

$$
\begin{align*}
& J^{(2)}(\mu) T_{s}-J^{(1)}(\mu)=\int_{-T}^{T}\left|f_{w}(t)\right|^{2} d t-\sum_{l=-M}^{M}\left|f^{T_{s}}\left(l T_{s}\right)\right|^{2} T_{s} \\
& \quad+2 \mu R^{(1)}(\mu)-2 \mu R^{(2)}(\mu) T_{s}+\sum_{l=-M}^{M}\left|g\left(l T_{s}\right)\right|^{2} T_{s} \\
& \quad-\int_{-T}^{T}|g(t)|^{2} d t . \tag{74}
\end{align*}
$$

Thus, we conclude that

$$
\lim _{T_{s} \rightarrow 0} J^{(2)}(\mu) T_{s}=J^{(1)}(\mu)
$$

Q.E.D.

Let $\mu_{R}^{(i)}, \mu_{\epsilon}^{(i)}, i=1,2,3$ be the parameters $\mu_{R}$ and $\mu_{\epsilon}$, respectively, for the three cases mentioned above. Then, based on Theorem 5, we claim that
i) $\lim _{N \rightarrow \infty} \mu_{R / \sqrt{T_{s}}}^{(3)}=\mu_{R}^{(2)} \sqrt{T_{s}}$
ii) $\lim _{N \rightarrow \infty} \mu_{\epsilon / \sqrt{T_{S}}}^{(3)}=\mu_{\epsilon / \sqrt{T_{S}}}^{(2)}$
iii) $\lim _{T_{s} \rightarrow 0} T_{s} \cdot \mu_{R / \sqrt{T_{s}}}^{(2)}=\mu_{R}^{(1)}$
iv) $\lim _{T_{s} \rightarrow 0} T_{s} \cdot \mu_{\epsilon / \sqrt{T_{s}}}^{(2)}=\mu_{\epsilon}^{(1)}$
where

$$
\begin{array}{ll}
R^{(3)}\left(\mu_{R}^{(3)} / \sqrt{T_{s}}\right)=R^{2} / T_{s}, & J^{(3)}\left(\mu_{\epsilon}^{(3)} \sqrt{T_{s}}\right)=\epsilon^{2} / T_{s} \\
R^{(2)}\left(\mu_{R / \sqrt{T_{s}}}^{(2)}\right)=R^{2} / T_{s}, & J^{(2)}\left(\mu_{\varepsilon}^{(2)} \sqrt{T_{s}}\right)=\epsilon^{2} / T_{s} \\
R^{(1)}\left(\mu_{R}^{(1)}\right)=R^{2}, & J^{(1)}\left(\mu_{\epsilon}^{(1)}\right)=\epsilon^{2} .
\end{array}
$$

Therefore, $T_{s} \cdot \mu_{R}^{(3)} / \sqrt{T_{s}}$ and $T_{s} \cdot \mu_{\epsilon}^{(3)} / \sqrt{T_{s}}$ are the approximations of the parameters $\mu_{R}^{(1)}$ and $\mu_{\epsilon}^{(1)}$ when $T_{s}$ is sufficiently small and $N$ is sufficiently large.
In conclusion, we have proved that for $N \rightarrow \infty$ and $T_{s} \rightarrow 0$ the solution $f_{*}^{N,} T_{s}\left(m T_{s}\right)$ of the numerical implementation of the extrapolating algorithms will tend towards the analog solution $f_{*}(t)$. However, the important question of how large (small) should we choose $N,\left(T_{s}\right)$ has not been addressed here.

## VI. Numerical Example

We illustrate the method with a numerical example by solving Problem 1 for periodic band-limited sequences. The performed iteration has the following steps.

Step 1: Set

$$
\mu_{L}=0, \mu_{U}=\sum_{l=-M}^{M}|g[l]|^{2} / 2 R^{2}
$$

Step 2: Set

$$
\mu=\frac{1}{2}\left(\mu_{L}+\mu_{U}\right)
$$

$f_{0}[l]=0$,

$$
\alpha=1 / 1+\mu
$$

Step 3: Form

$$
w_{n}[l]=(1-\alpha \mu) f_{n-1}[l]+\alpha\left(g[l]-f_{n-1}[l]\right) P_{M}[l]
$$

Step 4: Take the DFS $W_{n}[m]$ of $w_{n}[l]$.
Step 5: Form

$$
F_{n}[m]= \begin{cases}W_{n}[m] & |m| \leqslant K \\ 0 & |m|>K\end{cases}
$$

Step 6: Take the inverse $\operatorname{DFS} f_{n}[l]$ of $F_{n}[m]$.
Step 7: Evaluate

$$
R_{n}^{2}=\sum_{l=0}^{N-1}\left|f_{n}[l]\right|^{2} \quad L_{n}=\left[1-(1+\mu)^{-n}\right]^{2} R^{2}
$$

Then, if
a) $R_{n}^{2}>R^{2}$, set $\mu_{U}=\mu$ and go to step 2 .
b) $R_{n}^{2}<L_{n}$, set $\mu_{L}=\mu$ and go to step 2 .
c) $R^{2}(1-C)<R_{n}^{2}$, stop.
d) Otherwise, go to step 3 .

In the example below, the computations were carried out with a DFS of size $N=256$ and the various constants were set as follows:

$$
M=20 \quad K=15 \quad C=10^{-5}
$$



Fig. 2. $G[m]$ : The DFS of the data.


Fig. 3. (a) The given data $g[l]$ for $|l| \leqslant 20$. (b) The inverse $f[l]$ of the $K$-band-limited part of $G[m]$.


Fig. 4. The extrapolated $f_{*}[l]$ for (a) $\mu=0, n=500$. (b) $R^{2}=16.0$. (c) $R^{2} \stackrel{256 / 31 . ~(d) ~}{ } R^{2}=4.0$.

The data $g[l]$ were obtained from

$$
G[m]= \begin{cases}\frac{1}{31} & |m| \leqslant K \\ \text { white noise } & \text { otherwise. }\end{cases}
$$



Fig. 5. The DFS $F_{*}[m]$ of $f_{*}[l]$ for (a) $\mu=0, n=500$. (b) $R^{2}=16.0$. (c) $R^{2}=256 / 31$. (d) $R^{2}=4.0$.


Fig. 6. $R_{n}^{2}$ for various values of $\mu$ versus $n$.
$G[m]$ is shown in Fig. 2 and the given data $g[l]$ in Fig. 3(a). The unknown inverse $f[l]$ of the $K$-band-limited part of $G[m]$ is shown in Fig. 3(b). Its energy is $256 / 31$. The extrapolation was performed for four different cases $\mu=0, R^{2}=16.0, R^{2}=$ $256 / 31, R^{2}=4.0$. The results are shown in Fig. 4, and their Fourier transforms in Fig. 5. It appears that $R^{2}=4.0$ gives us the best solution. This is because the solution $f_{*}[l]$ only can


Fig. 7. The functions $R(\mu)$ and $J(\mu)$ for the presented example.
guarantee $\left\|f_{*}[l]-g[l]\right\|_{-M, M}^{2}$ is minimum over the family of $K$-band-limited $N$-periodic sequences whose energy is equal to or less than $R^{2}$. This does not imply that $\left\|f_{*}[l]-f[l]\right\|_{0, N-1}^{2}$ is minimum when $R^{2}=\|f[l]\|_{0, N-1}^{2}$. The case $\mu=0$ is the algorithm by Papoulis and Gerchberg. In Fig. $6, R_{n}^{2}$ is shown for the four cases. It is clear that for $\mu=0$ we have a diverging algorithm. Finally, in Fig. 7 the $R(\mu)$ and $J(\mu)$ of the given example are plotted as a function of $\mu$. We obtain similar behavior from Problems 2 and 3.

## VII. Conclusion

The applicability of the extrapolating algorithm proposed by Papoulis and Gerchberg was extended over noisy data by regularization. This was achieved by imposing energy constraints on the unknown band-limited signal or the noise. The discrete versions of the revised algorithm as well as the transition from the continuous algorithm to its digital implementation are presented. It is of interest to notice that by overcoming some difficulties on mathematics a similar algorithm can be derived if we consider $n(t)$ to be a stochastic process and seek for maximum likelihood solution. This result will be discussed in a future work.

## Appendix A

## The Prolate Spheroidal Functions

In this Appendix we simply state the definition of the prolate spheroidal function and the properties we need for the derivations in this work. For more details, see [1] or [11].
Definition: The prolate spheroidal functions are the eigenfunctions $\phi_{k}(t)$ of the equation

$$
\begin{equation*}
\int_{-T}^{T} \phi_{k}(\tau) \sin \sigma(t-\tau) / \pi(t-\tau) d \tau=\lambda_{k} \phi_{k}(t) \tag{A1}
\end{equation*}
$$

## Properties:

a) $\phi_{k}(t)$ are $\sigma$-band-limited, i.e.,

$$
\begin{equation*}
\phi_{k}(t) * \frac{\sin \sigma t}{\pi t}=\phi_{k}(t) \tag{A2}
\end{equation*}
$$

b) $\phi_{k}(t)$ are orthonormal in $(-\infty, \infty)$ and orthogonal in $(-T, T)$, i.e.,

$$
\begin{align*}
& \int_{-\infty}^{\infty} \phi_{k}(t) \phi_{l}(t) d t= \begin{cases}1 & k=l \\
0 & k \neq l\end{cases}  \tag{A3}\\
& \int_{-T}^{T} \phi_{k}(t) \phi_{l}(t) d t= \begin{cases}\lambda_{k} & k=l \\
0 & k \neq l\end{cases} \tag{A4}
\end{align*}
$$

c) The eigenvalues $\lambda_{k}$ are such that

$$
\begin{align*}
& 1>\lambda_{0}>\lambda_{1}>\cdots>\lambda_{k}>\cdots>0 \\
& \quad \text { and } \quad \lim _{k \rightarrow \infty} \lambda_{k}=0 \tag{A5}
\end{align*}
$$

d) Any $\sigma$-band-limited function $f(t)$ can be expressed as

$$
\begin{align*}
f(t) & =\sum_{k=0}^{\infty} a_{k} \phi_{k}(t) \\
a_{k} & =\int_{-\infty}^{\infty} f(t) \phi_{k}(t) d t \tag{A6}
\end{align*}
$$

and

$$
\begin{equation*}
\|f(t)\|^{2}=\sum_{k=0}^{\infty} a_{k}^{2} \tag{A7}
\end{equation*}
$$

e) Any function $g(t)$ with

$$
\|g(t)\|_{T}^{2}=\int_{-T}^{T}|g(t)|^{2} d t<\infty
$$

can be expressed in the interval $(-T, T)$ as

$$
\begin{align*}
g(t) & =\sum_{k=0}^{\infty} b_{k} \phi_{k}(t) \quad|t|<T \\
b_{k} & =\frac{1}{\lambda_{k}} \int_{-T}^{T} g(t) \phi_{k}(t) d t \tag{A8}
\end{align*}
$$

and

$$
\begin{equation*}
\|g(t)\|_{T}^{2}=\sum_{k=0}^{\infty} \lambda_{k} b_{k}^{2} \tag{A9}
\end{equation*}
$$

## Appendix B

The Discrete Prolate Spheroidal Sequences (DPSS)
Definition: The discrete prolate spheroidal sequences are the eigenvectors $\phi_{k}[\mathrm{~m}]$ of the equation

$$
\begin{align*}
& \sum_{n=-M}^{M} \frac{\sin \sigma(m-n)}{\pi(m-n)} \phi_{k}[n]=\lambda_{k} \phi_{k}[m] \\
& k=0, \cdots, 2 M \tag{B1}
\end{align*}
$$

## Properties:

a) $\phi_{k}[m]$ are $\sigma$-band-limited, i.e.,
$\phi_{k}[m] * \frac{\sin \sigma m}{\pi m}=\phi_{k}[m]$.
b) $\phi_{k}[m]$ are doubly orthogonal, i.e.,

$$
\begin{align*}
& \sum_{m=-\infty}^{\infty} \phi_{k}[m] \phi_{l}[m]= \begin{cases}1 & k=l \\
0 & k \neq l\end{cases}  \tag{B3}\\
& \sum_{m=-M}^{M} \phi_{k}[m] \phi_{l}[m]= \begin{cases}\lambda_{k} & k=l \\
0 & k \neq l .\end{cases} \tag{B4}
\end{align*}
$$

c) The eigenvalues $\lambda_{k}$ are such that
$1>\lambda_{0}>\lambda_{1}>\cdots>\lambda_{2 M}>0$.
d) There exist $\sigma$-band-limited sequences $\phi_{k}[m], k=$ $2 M+1, \cdots$, such that

$$
\begin{aligned}
& \sum_{m=-\infty}^{\infty} \phi_{k}[m] \phi_{l}[m]= \begin{cases}1 & k=l \\
0 & k \neq l\end{cases} \\
& k, l=0,1,2, \cdots
\end{aligned}
$$

and any $\sigma$-band-limited sequence $f[m]$ can be expressed as

$$
\begin{align*}
f[m] & =\sum_{k=0}^{\infty} a_{k} \phi_{k}[m] \\
a_{k} & =\sum_{m=-\infty}^{\infty} f[m] \phi_{k}[m] \tag{B6}
\end{align*}
$$

and

$$
\begin{equation*}
\|f[m]\|^{2}=\sum_{k=0}^{\infty} a_{k}^{2} \tag{B7}
\end{equation*}
$$

e) Any vector $(g[-M], \cdots, g[M])$ can be expressed as

$$
\begin{align*}
g[m] & =\sum_{k=0}^{2 M} b_{k} \phi_{k}[m] \quad|m| \leqslant M \\
b_{k} & =\frac{1}{\lambda_{k}} \sum_{m=-M}^{M} g[m] \phi_{k}[m] \quad k=0, \cdots, 2 M \tag{B8}
\end{align*}
$$

and

$$
\begin{equation*}
\|g[m]\|_{-M, M}^{2}=\sum_{k=0}^{2 M} \lambda_{k} b_{k}^{2} \tag{B9}
\end{equation*}
$$

It is clear that if $g[m]=f[m]$ for $|m| \leqslant M, b_{k}=a_{k}$ for $k=$ $0,1, \cdots, 2 M$.

## Appendix C

## The Periodic Discrete Prolate Spheroidal SEQuences (P-DPSS)

Definition: The periodic discrete prolate spheroidal sequences are the eigenvectors $\phi_{i}[m]$ of the equation, i.e.,

$$
\begin{align*}
& \sum_{n=-M}^{M} \frac{\sin \frac{\pi}{N}(2 K+1)(m-n)}{N \sin \frac{\pi}{N}(m-n)} \phi_{i}[n]=\lambda_{i} \phi_{i}[m] \\
& \quad i=0, \cdots, 2 M \tag{C1}
\end{align*}
$$

where

$$
2 K+1<N, \quad 2 M+1<N .
$$

## Properties:

a) The eigenvalues $\lambda_{i}$ are such that

$$
\begin{equation*}
1 \geqslant \lambda_{0} \geqslant \lambda_{1} \geqslant \cdots \geqslant \lambda_{2 M}>0 \quad \text { if } \quad M \leqslant K \tag{C2}
\end{equation*}
$$

or

$$
\begin{align*}
& 1 \geqslant \lambda_{0} \geqslant \lambda_{1} \geqslant \cdots \geqslant \lambda_{2 K}>\lambda_{2 K+1}=\cdots=\lambda_{2 M}=0 \\
& \quad \text { if } M>K . \tag{C3}
\end{align*}
$$

b) $\phi_{i}[m]$ are doubly orthogonal, i.e.,

$$
\begin{align*}
& \sum_{m=l_{0}}^{l_{0}+N-1} \phi_{i}[m] \phi_{j}[m]= \begin{cases}1 & i=j \\
0 & i \neq j\end{cases}  \tag{C4}\\
& \sum_{m=-M}^{M} \quad \phi_{i}[m] \phi_{j}[m]= \begin{cases}\lambda_{i} & i=j \\
0 & i \neq j\end{cases} \tag{C5}
\end{align*}
$$

where

$$
\begin{equation*}
i, j=0, \cdots, \min (2 K, 2 M) \tag{C6}
\end{equation*}
$$

c) $\phi_{i}[m], i=0, \cdots, \min (2 K, 2 M)$, are $K$-band-limited, i.e.,

$$
\begin{equation*}
\sum_{n=l_{0}}^{l_{0}+N-1} \frac{\sin \frac{\pi}{N}(2 K+1)(m-n)}{N \sin \frac{\pi}{N}(m-n)} \phi_{i}[n]=\phi_{i}[m] \tag{C7}
\end{equation*}
$$

d) When $M \leqslant K$, there exist $K$-band-limited $N$-periodic sequences $\phi_{i}[m], i=2 M+1, \cdots, 2 K$, such that

$$
\sum_{m=l_{0}}^{l_{0}+N-1} \phi_{i}[m] \phi_{j}[m]= \begin{cases}1 & i=j \\ \vdots & i \neq j\end{cases}
$$

and any $K$-band-limited $N$-periodic sequence $f[m]$ can be expressed as

$$
\begin{align*}
f[m] & =\sum_{i=0}^{2 K} a_{i} \phi_{i}[m] \\
a_{i} & =\sum_{m=l_{0}}^{l_{0}+N-1} f[m] \phi_{i}[m] \tag{C8}
\end{align*}
$$

and

$$
\begin{equation*}
\|f[m]\|_{l_{0}, l_{0}+N-1}^{2}=\sum_{i=0}^{2 K} a_{i}^{2} \tag{C9}
\end{equation*}
$$

e) When $M>K$, any $K$-band-limited $N$-periodic sequence $f[m$ ] can be expressed as

$$
\begin{align*}
f[m] & =\sum_{i=0}^{2 K} a_{i} \phi_{i}[m] \\
a_{i} & =\sum_{m=l_{0}}^{l_{0}+N-1} f[m] \phi_{i}[m] \tag{C10}
\end{align*}
$$

and

$$
\begin{equation*}
\|f[m]\|_{l_{0}, l_{0}+N-1}^{2}=\sum_{i=0}^{2 K} a_{i}^{2} \tag{C11}
\end{equation*}
$$

f) When $M \leqslant K$, any vector ( $g[-M], \cdots, g[M]$ ) can be expressed as

$$
\begin{align*}
g[m] & =\sum_{i=0}^{2 M} b_{i} \phi_{i}[m] \quad|m| \leqslant M \\
b_{i} & =\frac{1}{\lambda_{i}} \sum_{m=-M}^{M} g[m] \phi_{i}[m] \tag{C12}
\end{align*}
$$

and

$$
\begin{equation*}
\|g[m]\|_{-M, M}^{2}=\sum_{i=0}^{2 M} \lambda_{i} b_{i}^{2} \tag{C13}
\end{equation*}
$$

g) When $M>K$, any vector ( $g[-M], \cdots, g[M]$ ) can be expressed as

$$
\begin{align*}
& g[m]=\sum_{i=0}^{2 K} b_{i} \phi_{i}[m]+\tilde{g}[m] \quad|m| \leqslant M \\
& \sum_{m=-M}^{M} \tilde{g}[m] \phi_{i}[m]=0 \quad i=0, \cdots, 2 K  \tag{C14}\\
& b_{i}=\frac{1}{\lambda_{i}} \sum_{m=-M}^{M} g[\tilde{m}] \phi_{i}[m] \tag{C15}
\end{align*}
$$

and

$$
\begin{equation*}
\|g[m]\|_{-M, M}^{2}=\sum_{i=0}^{2 K} \lambda_{i} b_{i}^{2}+\|\tilde{g}[m]\|_{-M, M} \tag{C16}
\end{equation*}
$$

## Acknowledgment

The authors would like to thank Prof. A. Papoulis for his valuable comments.

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Wen Yuan Xu graduated from the Department of Mathematics, Nankai University, Tianjin, People's Republic of China, in 1963.
From 1963 to 1979 he taught at Nankia University. He has been working as a Research Associate at the Institute of Systems Science, Academia Sinica, Beijing, People's Republic of China, since 1979. From 1981 to 1983 he was a Visiting Scholar at the Polytechnic Institute of New York, Farmingdale. His research interests are in the areas of signal processing, information theory, and system identification.


Christodoulos Chamzas was born in Komotini, Greece, in 1951. He received the Dipl. Eng. degree in electrical and mechanical engineering from the National Technical University of Athens, Athens, Greece, in 1974 and the M.S. and Ph.D. degrees in electrical engineering in 1975 and 1979 from the Polytechnic Institute of New York, Farmingdale.
Since 1979 he has been an Assistant Professor with the Department of Electrical Engineering, Polytechnic Institute of New York. His primary interests are in digital signal processing and communication systems. He has held summer positions in Greece, England, and Portugal, as well as at Bell Laboratories, Holmdel, NJ.

Dr. Chamzas is a member of the Technical Chamber of Greece and Sigma Xi. He was the cowinner with A. Papoulis of the RADC Spectral Estimation Competition in 1978.


[^0]:    Manuscript received January 18, 1982; revised October 29, 1982.
    W. Y. Xu was on leave at the Department of Electrical Engineerng Polytechnic Institute of New York, Farmingdale, NY 11735. He is with the Institute of Systems Science, Academia Sinica, Beijing, People's Republic of China.
    C. Chamzas is with the Department of Electrical Engineering, Polytechnic Institute of New York, Farmingdale, NY 11735.
    ${ }^{1}$ The band-limited extrapolation algorithm was first presented by A. Papoulis in 1973 in an internal JSTAC report [2]; it was published independently by R. W. Gerchberg in 1974 [3]; the proof of the convergence was given by A. Papoulis in 1975 [4].

